FOR RELEASE: 06/23/11: CIA-RDP86-00513R001033800004-6

TYSAROWSKI, W.; KWIEK, S.; MIGDALSKA, B.

Behavior of hemoproteins in the presence of acid-fast bacilli. II. Effect of oxy- and methemoglobin on respiration of Mycobacterium phlei and Mycobacterium tuberculosis strain C. Acta microb. polon 5 no.1-2:65-68 1956.

1. Z Zakladu Biochemii i Mikrobiologii Instytutu Gruzlicy w Warszawie.

(HEMOGLOBIN, effects,

oxy- & methemoglobin on M. phlei & M. tuberc. resp. (Pol))

(MYCOBACTERIUM, effect of drugs on,

phlei, oxy-& methemoglobin (Pol)) (MYCOBACTERIUM TUBERCULOSIS, effect of drugs on,

oxy- & methemoglobin (Pol))

. P. C.

TYSAROWSKI, Wieslaw; KWIEK, Stanislaw; MIGDALSKA, Barbara

The behavior of heomproteins in the presence of acid resistant bacteria. I. Reduction of methemoglobin in clutures of Tubercle bacilli. Gruzlica 23 no.1:13-20 Jan '55.

CIA-RDP86-00513R001033800004

1. Z Zakladu Biochemii. Kierownik: prof.dr G. Bagdasarian i Zakladu Mikrobiologii Kierownik: doc.dr M. Buraszewska. Institutu Gruzlicy, Dyrektor: prof.dr J. Misiewicz. Warszawa, Plocka 26. (MYCOBACTERIUM TUBERCULOSIS, culture

methemoglobin reduction to hemoglobin)

methemoglobin reduction in M. tuberc.culture (HEMOGLOBIN

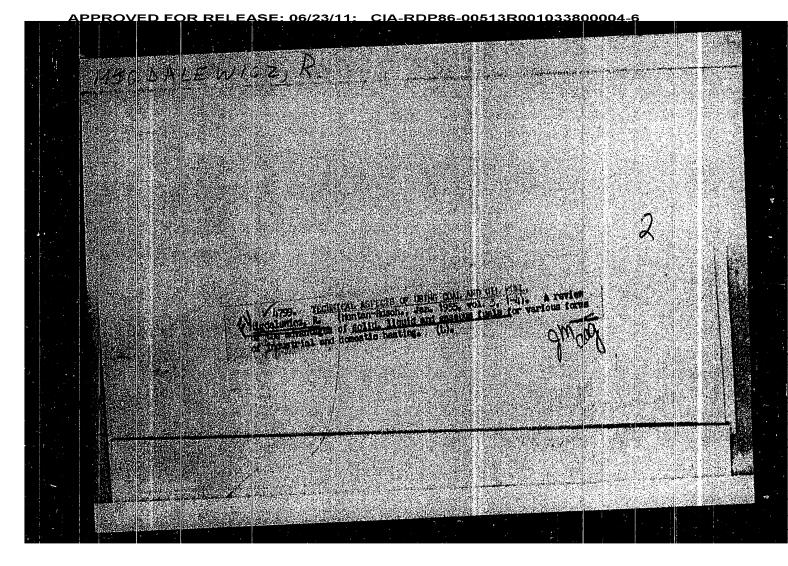
NICONIONICH, F. A.

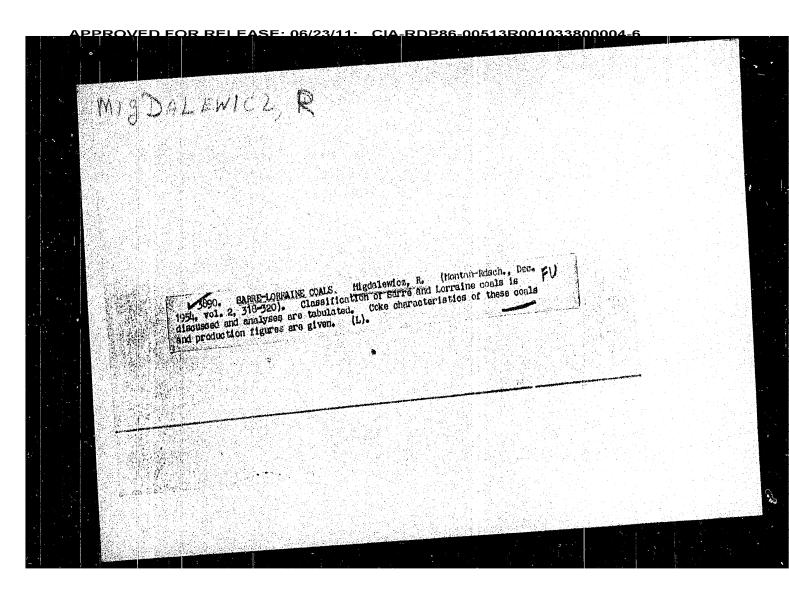
36915. Streptolokkovova cerekrostneva infoktsiya ne skorlatinnykh otdelenimakh. Voprosy pediatrii i o'dararu meteriastva lijetstva. 1850, vm. 5, 5. 25-27

MIGDALOVICH, F. A., KURILOVA, O. M., I HAKHLIHA, I. L.

SO: Letopis' Americal rock States, So. NO, 1989

GEFTER, S.P.; MIGDALOVICH, B.M.; GOLUBEVA, T.S. Skin tuberculin sensitivity in pulmonary tuberculosis during antibacterial therapy. Probl. tub. 41. no.3:34-37'63. (MIRA 16:9) 1. Iz kafedry legochnogo tuberkuleza (zav. - prof. A. Ya. TSigel'nik) I Leningradskogo meditsinskogo instituta imeni akademika I.P.Pavlova. (TUBERCULIN-TESTING) (CHEMOTHERAPY)





particle 1-forbidden transitions in which the orbital angular momentum of the quasi-particle changes by two units. "The authors thank Yu. V. Gaponov and E. Ye. Sapershteyn for valuable discussions." Orig. art, has: 18 formulas and 1 table. L 2232-66 ACCESSION NR: AP5020250 ASSOCIATION: Moskovskiy inzhenerno-fizicheskiy institut (Moscow Engineering-SUB CODE: NP Physics Institute) 44,55 ENCL: 00 SUBMITTED: 29Dec64 OTHER: NR REF SOV: 005

L 2232-66 EWT (m) DIAAP

ACCESSION NR: AP5020250

UR/0367/65/002/001/0028/0034

AUTHOR: Migdal, A.

TITLE: Beta decay in nuclei

SOURCE: Yadernaya fizika, v. 2, no. 1, 1965, 28-34

TOPIC TAGS: Beta decay, particle interaction, nuclear spin, forbidden transition

ABSTRACT: The method of interacting quasiparticles, developed by one of the authors (Migdal, Nucl. Phys. v. 57, 29, 1964), is used to analyze beta decay in nuclei. The probabilities for the allowed beta transitions are calculated, with account taken of the interaction between quasiparticles, by calculating the matrix elements for the Fermi and Gamow-Teller transitions. It is shown that the Fermi matrix elements can be calculated accurately without taking Coulomb interaction into account. In the case of Gamow-Teller transitions in mirror nuclei, the field satisfies an equation identical with that for the polarizability of the daughter nucleus in the field. The presence of a spin-spin interaction between quasiparticles in Gamow-Teller transitions and to the appearance of the group of single-

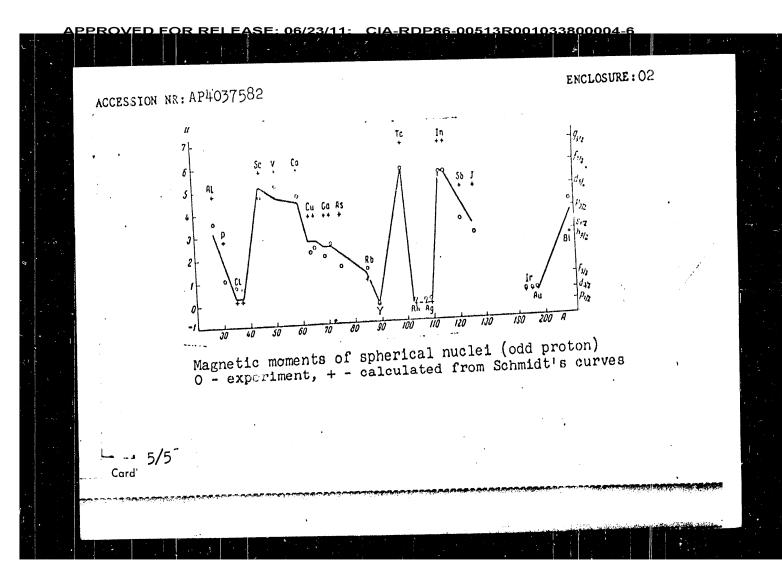
1/2 Card

MIGDAL, Arkadiy Deynusovich (1911-); DUBNOVA, V.Ya., red. (Theory of finite Fermi systems and the properties of atomic nuclei] Teorifa konechnykh fermi-sistem i svoistv atomrykh iader. Moskva, Nauka, 1965. 572 p. (MIRA 18:12) 1. Chlen-korrespondent AN SSSR (for Migdal).

MIGDAL, Arktdly Benediktovich; FAYNBOYM, I.B., red.

[Modern approach to maclear theory] Sevrementyl podkhod k teoril ladra. Moskva, Zmanie, 1965. 42 p. (Novoe v zhizni, nauke, tekhnike. IX Seriia: Fizika, matematika, astronomila, no.14)

1. Chlen-korrespondent AN SUSR (for Migdal).



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ACCESSION NR: AP4037582

ASSOCIATION: None

SUBMITTED: 01Aug63 DATE ACQ: 09Jun64

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OTHER: 005

SUB CODE: GP, NP

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APPROVED FOR RELEASE: 06/23/11: CIA-RDP86-00513R001033800004-6

ACCESSION NR: AP4037582

tion, the spin-orbit interaction gives rise to terms proportional to the orbital angular momentum of the odd particle, which is shown to make in some cases an appreciable contribution to the magnetic moment. A simple expression is derived for the sum of the magnetic moments of the neutron and the proton. The magnetic moments of the spherical nuclei obtained from the expression derived in this work agree well with experiment and with calculations made by others. The spin-spin interaction constant for the magnetic moments of deformed nuclei are found to coincide with those obtained from magnetic moments of spherical nuclei. The equations obtained can be used to calculate the magnetic form factors of the nuclei, and in particular, to determine the nuclear magnetic multipoles. "The author is grateful to A. A. Lushnikov, E. E. Sapershteyn, M. A. Troitskiy, and V. A. Khodel' for interesting discussion, and to V. N. Guman for compiling the table of integrals of the radial functions." Orig. art. has: 2 figures, 67 formulas, and 1 table.

Card 2/5

APPROVED FOR RELEASE: 06/23/11: CIA-RDP86-00513R001033800004-6

ACCESSION NR: AP4037582

s/0056/64/046/005/1680/1699

AUTHOR: Migdal, A. B.

TITLE: Magnetic moments of nuclei

SOURCE: Zh. eksper. i teor. fiz., v. 46, no. 5, 1964, 1680-1699

TOPIC TAGS: nuclear magnetic moment, Fermi surface, spin orbit coupling, spin spin coupling, deformed nucleus

ABSTRACT: As a further development of his quantitative approach to the calculation of nuclear phenomena (ZhETF v. 43, 1940, 1962; v. 44, 1703, 1963; v. 45, 1036, 1963), based on the introduction of constants which characterize the properties of nuclear matter and which are the same for all nuclei and for all types of transitions near the Fermi boundary, the author derives expressions for the magnetic moments of nuclei in which exact account is taken for the interaction between the nucleons. The magnetic moments are corrected for spin-orbit interaction with a high degree of accuracy. In addi-

Card --/

MIGDAL, A. 3. "Development of the Phenomenological Theory of the Nucleus." report submitted for All-Union Conf on Nuclear Spectroscopy, Tbilisi,  $1^{\frac{1}{4}-22}$ Feb 64.

MIGDAL, A.B.; LARKIN, A.I. Phenomenological approach to the theory of the nucleus. Zhur. eksp. i teor. fiz. 45 no.4:1036-1050 0 163. (MIRA 16:11)

L 10193-63 ACCESSION NR: AP3000071

0

find the probabilities of electromagnetic transitions in nuclei. Orig. art. has: 63 formulas, of which 9 are in diagram form.

ASSOCIATION: none

SUBMITTED: 25Dec62

DATE ACQ: 12Jun63

ENCL:

SUB CODE: PH

NR REF SOV: 007

OTHER: 000

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<u> APPROVED FOR RELEASE: 06/23/11: CIA-RDP86-00513R001033800004-6</u>

L 10193-63

EWT(1)/BDS-AFFTC/ASD-GO

ACCESSION NR: AP3000071

8/0056/63/044/005/1703/1718

AUTHOR: Larkin, A. I.; Migdal, A. B.

TITLE: Theory of a superfluid Fermi liquid. Application to the nucleus.

SOURCE: Zhurnal eksper. i teoret. fiziki, v. 44, no. 5, 1963, 1703-1718

TOPIC TAGS: Superfluid Fermi liquid, strong-interactions

ABSTRACT: A method is given for treating systems of strongly interacting particles, in which the observable quantities are expressed in terms of several constants that are introduced into the theory somewhat in the way that the masses and charges of particles are introduced in the theory of dispersion relations. The two-particle spectrum and the reaction of the system to an external field are determined by the two-particle Green's function, and it is shown that the finding of these quantities reduces to the solution of a simple equation analogous to the Schroedinger equation for two interacting particles in a potential well. The Landau theory of the Fermi liquid is extended to the case of superfluidity for systems of finite size. An equation is obtained which makes it possible to

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<u> APPROVED FOR RELEASE: 06/23/11: CIA-RDP86-00513R001033800004-6</u>

ACCESSION NR: AT4014034

Cooper pairing and the finite dimensions of the system. The latter is taken into account readily by the Green's function method. To describe the behavior of particles in an external field, a system of equations is formulated, the solution of which will yield the exact intensity of single-particle transitions and the magnetic moments of the nuclei. The coefficients of the solutions are determined by experiments. Further expansion of the theory to collective transitions is proposed. Orig. art. has: 22 formulas and 1 table.

ASSOCIATION: Fizicheskiy institut AN ArmSSR (Physics Institute, AN ArmSSR)

SUBMITTED: 00

DATE ACQ: 20Feb64

ENCL: 00

SUB CODE: PH

NO REF SOV: 004

OTHER: 002

Card 2/2

APPROVED FOR RELEASE: 06/23/11: CIA-RDP86-00513R001033800004-6

ACCESSION NR: AT4014034

S/2918/63/000/000/0363/0372

AUTHOR: Migdal, A. B.

TITLE: Phenomenological approach to nuclear theory

SOURCE: AN ArmSSR. Fizicheskiy institut. Voprosy\* fiziki elementarny\*kh chastits, 1963, 363-372

TOPIC TAGS: nuclear theory, quantitative nuclear theory, Fermi liquid theory, Cooper pair, Green's function, particle system in field, single particle transition, magnetic moment, collective transition

ABSTRACT: In an attempt to construct a quantitative nuclear theory, the author proposes to extend to the case of the nucleus the program used by Landau in the Fermi liquid theory (ZhETF v. 30, 1958, and v. 32, 59, 1957) for a system consisting of one type of particles, to include a system of two types of particles, with allowance for the

Card 1/2

Theory of a Fermi...

S/056/62/043/005/048/058 B125/B104

The dependence of the energy density wof the system on the densities  $\boldsymbol{n}_{\underline{a}}$  and  $n_{\hat{b}}$  of particles of the types a and b is given by

$$W = \frac{1}{2} K \frac{[(n_a + n_b) - 2n_0]^2}{2n_0} + \beta \frac{(n_a - n_b)^2}{2n_0}.$$
 (54).

 $n_{_{\mbox{\scriptsize O}}}$  is the equilibrium density at  $n_{_{\mbox{\scriptsize B}}}=n_{_{\mbox{\scriptsize D}}}$  when external fields are absent. In the case of nuclei, the second member on the right side of Eq. (54) goes over into the term  $\beta (N-Z)^2/\Lambda$  of the Weitzsäcker equation for the nuclear mass deficiency. The formulas

$$1 - \{f_{aa}^{k} + f_{ab}^{k}\} \equiv 1 - \chi_{0}^{k} = (vp_{0}/n)_{a}n_{0}/3K,$$

$$1 - \{f_{aa}^{k} - f_{ab}^{k}\} \equiv 1 - \eta_{0}^{k} = (vp_{0}/n)_{a}n_{0}/6\beta.$$
(59)

interrelate the scattering amplitudes and the rigidity coefficients K and B. The approximation of pair collisions in nuclei proved to be unusable. There are 4 figures. SUBMITTED:

June 21, 1962 Card 4/4

Theory of a Fermi...

S/056/62/043/005/048/058 B125/B104

(14).

$$f = f^{\omega} + \int f^{\omega} \frac{v \, \mathbf{k} \, \mathbf{n}_1}{\omega - v \, \mathbf{k} \, \mathbf{n}_1} f \frac{do_1}{4\pi},$$

$$g = g^{\omega} + \int g^{\omega} \frac{v \, \mathbf{k} \, \mathbf{n}_1}{\omega - v \, \mathbf{k} \, \mathbf{n}_1} g \frac{do_1}{4\pi},$$
(11)

 $\overrightarrow{n}$  and  $\overrightarrow{n}'$  are the unit vectors in the directions of  $\overrightarrow{p}$  and  $\overrightarrow{p}'$ . The density of the free particles is  $n = 2(4\pi/3)p_0^3/(2\pi)^3$ , also for a system consisting of two types of particles. The effective mass is

$$m_a^* = 1 + \frac{1}{3} \{ f_{aa}^\omega + f_{ab}^\omega \}_1 = 1 + \frac{1}{3} (\chi_1^\omega)_a.$$
 (53).

The quantities  $\chi$  = f ab and  $\gamma$  = f satisfy equations of the type

$$\chi = \chi^{\omega} + \int \chi^{\omega} \frac{v k n_1}{\omega - v k n_1} \chi^{\frac{do_1}{4\pi}},$$

 $\chi = \chi^k + \int \chi^k \frac{\omega}{\omega - v k n_1} \chi^{\frac{do_1}{4\pi}}.$ 

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APPROVED FOR RELEASE: 06/23/11: CIA-RDP86-00513R001033800004-6

Theory of a Fermi...

3/056/62/043/005/048/058 B125/B104

$$\Gamma_{-}(p, p'; q) = U(p, p'; q) - \frac{1}{2} \int U(p, p_1; q) G(p_1 + \frac{q}{2}) G(p_1 - \frac{q}{2}) \Gamma(p_1, p'; q) \frac{d^4p_1}{(2\pi)^4},$$
(1)

is then converted into the matrix equations

$$\Gamma_{aa}^{-} = \Gamma_{aa}^{\omega} + \Gamma_{aa}^{\omega} A_{a} \Gamma_{aa} + \Gamma_{ab}^{\omega} A_{b} \Gamma_{ba}.$$

$$\Gamma_{ba} = \Gamma_{ba}^{\omega} + \Gamma_{ba}^{\omega} A_{a} \Gamma_{aa} + \Gamma_{bb}^{\omega} A_{b} \Gamma_{ba}.$$
(71)

for the amplitudes  $l_{aa}$ ,  $l_{aa}^{\omega}$  of the scattering of equal particles and amplitudes  $l_{ab}^{\omega}$ ,  $l_{ab}^{\omega}$  of the scattering of the particle a on the particle b. The equation

$$F(\mathbf{n}, \sigma; n', \sigma') = f(\mathbf{n}, n') + g(\mathbf{n}, n') \sigma \sigma'$$
 (13)

relating the scattering amplitude as derived from (71) to the spin operators has an exchange character. Therefore,

Card 2/4

s/056/62/043/005/048/058 B125/B104 Theory of a Fermi liquid consisting of two types of particles. Migdal, A. B. AUTHOH: Application to the nucleus Zhurnal eksperimental noy i teoreticheskoy fiziki, v. 43, TITLET no. 5(11), 1962, 1940 - 1952 TEXT: The theory of L. D. Landau (ZhETF, 30, 1058, 1956; ZhETF, 35, 97, PERIODICAL: 1956) for calculating the basic properties of a homogeneous system, consisting of one type of strongly interacting Fermi particles, from the amplitude of the forward scattering of quasiparticles on the Fermi surface is extended to a system of two different types of particles. This generalized theory makes it possible to calculate the basic properties of the nuclei. The results of Landau's theory apply also to two types of particles if the scattering amplitude i, the complex U of all diagrams particles if the Scattering amplitude, the complex of all the limits having no parts linked by two horizontal lines; and the limits of the scattering amplitudes are two-rowed matrices in isotopic-spin space. Landau's equation Cará 1/4

PROVED FOR RELEASE: 06/23/11: CIA-RDP86-00513R001033800004-6

MICDAL A.B.

s/089/62/013/006/019/027 B102/B186

AUTHORS:

G. T. and M. R.

TITLE:

Nauchnaya konferentsiya Moskovskogo inzhenerno-fizicheskogo instituta (Scientific Conference of the Moscow Engineering Physics Institute) 1962

PERIODICAL: Atomnaya energiya, v. 13, no. 6, 1962, 603 - 606

TEXT: The annual conference took place in May 1962 with more than 400 delegates participating. A review is given of these lectures that are assumed to be of interest for the readers of Atomnaya energiya. They are following: A. I. Leypunskiy, future of fast reactors; A. A. Vasil'yev, design of accelerators for superhigh energies; I. Ya. Pomeranchuk, analyticity, unitarity, and asymptotic behavior of strong interactions at high energies; A. B. Migdal, phenomenological theory for the many-body problem; Yu. D. Fiveyskiy, deceleration of medium-energy antiprotons in matter; Yu. M. Kogan, Ya. A. Iosilevskiy, theory of the Mössbauer effect; M. I. Ryazanov, theory of ionization losses in nonhomogeneous medium; Yu. B. Ivanov, A. A. Rukhadze, h-f conductivity of subcritical plasma;

Card 1/m

S/056/61/040/002/045/C47 B102/B201

Single-particle ...

(43) 
$$A_{\lambda} = (e_{\lambda}' + E_{\lambda}')/(E_{\lambda} + E_{\lambda}' + \mu^{+} - \mu^{-}),$$

$$B_{\lambda} = (E_{\lambda} - e_{\lambda})/(E_{\lambda} + E_{\lambda}' + \mu^{+} - \mu^{-}),$$

$$C_{\lambda} = -\Delta (E_{\lambda} - e_{\lambda})/(E_{\lambda} + E_{\lambda}' + \mu^{+} - \mu^{-}) (E_{\lambda} - e_{\lambda}').$$

$$\Delta = \gamma_1 \sum_{\lambda} C_{\lambda},$$

$$C_{\lambda} = -\Delta (E_{\lambda} - \epsilon_{\lambda}) / (E_{\lambda} + E_{\lambda} + \mu^{+} - \mu^{-}) (E_{\lambda} - \epsilon_{\lambda})$$

$$1 = -\gamma_{1} \sum_{\lambda} \{ (E_{\lambda} - \epsilon_{\lambda}) / (E_{\lambda} + E_{\lambda} + \mu^{+} - \mu^{-}) (E_{\lambda} - \epsilon_{\lambda}) \}$$

$$1 = -\gamma_{1} \sum_{\lambda} \{ (E_{\lambda} - \epsilon_{\lambda}) / (E_{\lambda} + E_{\lambda} + \mu^{+} - \mu^{-}) (E_{\lambda} - \epsilon_{\lambda}) \}$$

 $W_{0}(N)$  and  $W_{s}(N)$  are ground-state and excited-state energies of the N-particle system, respectively. There are 9 figures and 8 Soviet-bloc references.

Moskovskiy inzhenerno-fizicheskiy institut (Moscow Institute of Physics and Engineering)

September 4, 1960 SUBMITTED:

Card 8/8

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s/056/61/040/002/045/047 3102/3201

Single-particle ...

(39) 
$$E_{1\lambda}^{\nu} = W_{s\nu}(N+1) - W_{0}(N) > W_{0}(N+1) - W_{0}(N) \equiv \mu_{1}^{+}, \\ E_{2\lambda}^{\nu} = W_{0}(N) - W_{s\nu}(N-1) < W_{0}(N) - W_{0}(N-1) \equiv \mu_{1}^{-},$$

(40) 
$$\pm \sqrt{\Lambda^{2} + (e_{\lambda} - \mu^{+})^{2}} > \mu_{1}^{+} - \mu^{+}, \quad \pm \sqrt{\Lambda^{2} + (e_{\lambda} - \mu^{-})^{2}} < \mu_{1}^{-} - \mu^{-}.$$
 (40)

(41) 
$$E_{1\lambda} = \mu^+ + V \overline{\Delta^2 + (\epsilon_{\lambda} - \mu^+)^2}, \quad E_{2\lambda} = \mu^- - V \overline{\Delta^2 + (\epsilon_{\lambda} - \mu^-)^2}$$
 (41)

(42) 
$$V^{\overline{\Delta^2 + (\epsilon_{\lambda} - \mu^{+})^2}} \equiv E_{\lambda}, \quad \epsilon_{\lambda} - \mu^{+} \to \epsilon_{\lambda};$$

$$V^{\overline{\Delta^2 + (\epsilon_{\lambda} - \mu^{-})^2}} \equiv E_{\lambda}, \quad \epsilon_{\lambda} - \mu^{-} \to \epsilon_{\lambda}.$$

card 7/8

s/056/61/040/002/045/047 B102/B201 Single-particle ... (111) is obtained instead of (11). Numerous relations are derived; among others  $(E_{1\lambda}^{\mathsf{v}} - e_{\lambda}) A_{\lambda}^{\mathsf{v}} = -\Delta C_{\lambda}^{\mathsf{v}}, \quad (E_{2\lambda}^{\mathsf{v}} - e_{\lambda}) B_{\lambda}^{\mathsf{v}} = \Delta D_{\lambda}^{\mathsf{v}},$ (36) $(E_{1\lambda}^{\nu} + \epsilon_{\lambda} - 2\mu^{+})C_{\lambda}^{\nu} = -\Delta A_{\lambda}^{\nu}, \quad (E_{2\lambda}^{\nu} + \epsilon_{\lambda} - 2\mu^{-})D_{\lambda}^{\nu} = \Delta B_{\lambda}^{0},$  $\Delta = \gamma_1 \sum_{\lambda} C_{\lambda}^{\nu}.$ (37)  $E_{1\lambda}^{0} = \mu^{+} \pm \sqrt{\Delta^{2} + (\epsilon_{\lambda} - \mu^{+})^{2}}, \quad E_{2\lambda}^{\vee} = \mu^{-} \pm \sqrt{\Delta^{2} + (\epsilon_{\lambda} - \mu^{-})^{2}}. \quad (37)$  $G_{\lambda}^{+}(\tau) = -i \sum_{s} |(a_{\lambda}^{+})_{s0}|^{2} \exp \{-i[W_{s}(N+1) - W_{0}(N)] \tau\},$ (38) $G_{\lambda}^{-}(\tau) = i \sum_{s} |(a_{\lambda})_{s0}|^{2} \exp \{i[W_{s}(N-1)-W_{0}(N)]\tau\},$ (38) Card 6/8 rge  $(a_{\lambda}^{+})_{s0} = (\Phi_{s} (N + 1), a_{\lambda}^{+} \Phi_{0} (N))$  is  $(a_{\lambda})_{s0} = (\Phi_{s} (N - 1), a_{\lambda} \Phi_{0} (N))$ :

<u> APPROVED FOR RELEASE: 06/23/11: CIA-RDP86-00513R001033800004-6</u>

Single-particle ...

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 $\Delta\,\text{are}$  discussed in Chapter 6. The graph equation

 $\Delta^* (x, p, \varepsilon) := (2\pi)^{-4} \int \Gamma^* (x - x', p, \varepsilon, p', \varepsilon_0) F(x, p', \varepsilon') dx' dp' d\varepsilon'.$ 

 $\Delta = \frac{1}{\sqrt{1 + \frac{1}{2}}}$  holds. By

one obtains

(20)

 $\Delta^{\bullet}(x) = (2\pi)^{-4} \int \gamma^{\bullet}(p') F(x, p', \epsilon') dp' d\epsilon',$   $\gamma^{\bullet}(p') = \frac{1}{4\pi} \int \Gamma^{\bullet}(x - x', p_0, \epsilon_0, p', \epsilon_0) dx' d\omega_{p'},$ 

for p = p<sub>0</sub> and  $\xi = \xi_0$ , if  $\Delta^{\#}(x) \equiv \Delta^{\#}(x, p_0, \xi_0)$ . The connection between the equations for the Green function and the vertex part is examined by the method of the two Green functions in Chapter 7. The second Green function

is defined by the method by L. P. Gor'kov.

(28)

 $(i\partial/\partial \tau - H) \widetilde{G}(x, x') = \delta(x - x') + i\widetilde{\Delta}(x) F(x, x'),$   $(i\partial/\partial \tau + H^{*} - 2\mu) \widetilde{F}(x, x') = -i\widetilde{\Delta}^{*}(x) \widetilde{G}(x, x'),$ 

 $\widetilde{\Delta}^{\bullet}(x) = \int \widetilde{\gamma}^{\bullet}(\rho') \widetilde{F}(x, \rho', e') d\rho' de' / (2\pi)^{4}.$ 

is obtained. The application of the system of equations to the nucleus is discussed in Chapter 8. For  $\tau < 0$ ,

Card 5/8

Single-particle ...

S/056/61/040/002/045/047 B102/B201

are obtained; if the renormalized Green function  $\widetilde{G}(G = (1 - \delta \Sigma_{R}/\sigma E_{0}))$  is developed according to the eigenvalues  $\psi_{\lambda}$ ,  $\widetilde{G}(\vec{r}, \vec{r}', E)$   $= \sum_{\lambda \lambda} G_{\lambda \lambda}, (E) \psi_{\lambda}(\vec{r}) \psi_{\lambda}^{*}, (\vec{r}'), \qquad (10)$ 

is obtained, and, without pairing,  $G_{\lambda\lambda}^{0}$ , =  $\delta_{\lambda\lambda}$ ,/( $\xi - \xi_{\lambda} + i\alpha | \xi - \xi_{\lambda_0}| (\xi - \xi_{\lambda_0})$ ), where  $\xi_{\lambda_0}$  is the highest filled level. The proper-energy part  $\xi_{k}$  which is connected with pair correlation is studied in Chapter 5.

 $\sum_{k}$  For a square-well potential U(r):

 $(\Sigma_{k})_{\lambda\lambda}$ , =  $\partial_{\lambda\lambda}$ ,  $|\Delta_{\lambda}|^{2}/(\varepsilon_{\lambda} + \varepsilon_{\lambda} - 2\mu)$ , while in all other cases  $((\Sigma_{k})_{\lambda\lambda}, = \sum_{\lambda} \Delta_{\lambda\lambda}, \Delta_{\lambda\lambda}, \Delta_{\lambda\lambda}/(\varepsilon + \varepsilon_{\lambda}, -2\mu))$ . The equations for the vertex part card 4/8

APPROVED FOR RELEASE: 06/23/11: CIA-RDP86-00513R001033800004-6

Single-particle ...

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The analytical properties of the proper-energy parts examined in Chapter 3.  $\Sigma = \Sigma_0 \left( \mathbf{r}, \, \mathbf{r}_1, \, \mathbf{e} \right) + \alpha \left( \mathbf{r}, \, \mathbf{r}_1 \right) | \mathbf{e} - \mathbf{e}_0 | \left( \mathbf{e} - \mathbf{e}_0 \right) | \mathbf{i} + 2\pi^{-1} \beta \left( \mathbf{r}, \, \mathbf{r}_1 \right) | \mathbf{e} - \mathbf{e}_0 |^3 \ln \left[ i \mathbf{e}_0 / (\mathbf{e} - \mathbf{e}_0) \right],$ 

is thus obtained, where  $\sum_0$  is a function of  $\epsilon$  which has no singularity at  $\epsilon \rightarrow \epsilon_0$  ( $\sum_0$  has a singularity at the distance  $\sim \epsilon_0$  from the real axis). Chapter 4 deals with the approximation of independent quasiparticles.

Gears with the approximate 
$$2\pi f$$
  $d\mathbf{r}_1 = (2\pi)^{-3} \sum_R (\mathbf{r}, p, \epsilon) G(\mathbf{p}, \mathbf{r}', \epsilon) e^{i\mathbf{p}\mathbf{r}} d\mathbf{p} \approx \int \Sigma_R(\mathbf{r}, | \mathbf{r} - \mathbf{r}_1|, \epsilon) G(\mathbf{r}_1, \mathbf{r}', \epsilon) d\mathbf{r}_1 = (2\pi)^{-3} \sum_R (\mathbf{r}, p, \epsilon) G(\mathbf{p}, \mathbf{r}', \epsilon) e^{i\mathbf{p}\mathbf{r}} d\mathbf{p} \approx (\Sigma_R(\mathbf{r}, p_0, \epsilon_0) + (\partial \Sigma_R(\mathbf{r}, p_0, \epsilon_0)/\partial \epsilon_0) (\epsilon - \epsilon_0) + (\partial M/p_0) (\partial \Sigma_R(\mathbf{r}, p_0, \epsilon_0)/\partial \epsilon_0) (\mathbf{p}^2/2M - p_0^2/2M)) G(\mathbf{r}, \mathbf{r}', \epsilon),$ 

$$\{\epsilon - \mathbf{p}^2/2M_{eff} - U(\mathbf{r})\} [1 - (\partial \Sigma_R(p_0, \epsilon_0)/\partial \epsilon_0)] \widetilde{G}(\mathbf{r}, \mathbf{r}', \epsilon) = \delta(\mathbf{r} - \mathbf{r}') + \int d\mathbf{r}_1 \Sigma_R(\mathbf{r}, \mathbf{r}_1, \epsilon) \widetilde{G}(\mathbf{r}, \mathbf{r}', \epsilon);$$

$$M_{eff} = M \frac{1 - \partial \Sigma_R/\partial \epsilon_0}{1 + (\partial \Sigma_R/\partial p_0) (M/p_0)},$$

$$Card 3/3 \qquad U(\mathbf{r}) = \frac{\Sigma_R(\mathbf{r}, p_0, \epsilon_0) - \epsilon_0 \partial \Sigma_R/\partial \epsilon_0 - (p_0, 2) (\partial \Sigma_R/\partial p_0)}{1 - \partial \Sigma_R/\partial \epsilon_0}.$$

APPROVED FOR RELEASE: 06/23/11: CIA-RDP86-00513R001033800004-6

<u> APPROVED FOR RELEASE: 06/23/11: \_CIA-RDP86-00513R001033800004-6</u>

Single particle ...

s/056/61/040/002/045/047 B102/B201

to consist of N particles of one kind; it is found, however, that two inde pendent systems of equations are obtained for a system of two kinds of Fermi particles (e.g., neutrons and protons in the nucleus); only the graphs as a whole become somewhat complicated. The application of results to real nucles is discussed by the example of the twofold magic nucleus (without pair correlation) Ph208. If the model of the free particles were correct, it would have to be possible to obtain the excited state of Pb with 2.6 Mev and the spin 3, which is regarded as consisting of quasiparticles and holes, from the ground states of  $Pb^{209}$  and  $Pb^{207}$ . In the actual fact, however, 3 5 Mev and a spin # 3 result instead of 2.6 Mev. The difficulties involved are designated as "three-spin problem"; its solution requires the introduction of an excitation interaction. The paper consists of eight chapters. Chapter ! (introduction) offers an exposition of the problems, the phenomenological bases of the method and a discussion of the accuracy. Chapter 2 contains the Dyson equation in coordinate representation for the singleparticle Green function and a brief discussion of the respective graph equation for the compact part of the proper energy

Card 2/8

\$/056/61/040/002/045/047 B162/3201

AUTHOR:

Migdal, A B.

TITLE:

Single-particle excitation and superfluidity in systems of Fermi

particles with arbitrary interaction. Application to the

nucleus

PERIODICAL:

Zhurnal eksperimental noy i teoreticheskoy fiziki, v 40.

no. 2, 1961, 684-697

TEXT: The author wanted to prove that a study of the analytical properties of Green functions when taking account of pair correlations and with any interaction of the particles of the system yields equations describing the single-particle excitation, if the excitation energy is small compared with the chemical potential of the system; these equations accurately describe the excited states of a system consisting of a finite number (N) of particles including terms of the order of N-1/3. One must know the spectrum of singleparticle excitations without pair correlation in order to be able to determine the energy of the ground state and of the first excited single-particle states from the system of equations. To begin with, the system is assumed

Card 1/8

Superfluidity and Moments of Inertia of Nuclei

SOV/56-37-1-38/64

a method of calculating the moment of inertia, calculation of the moment of inertia for the oscillator potential, calculation of the moment of inertia for a rectangular potential well, moments of inertia corresponding to neutrons and protons, gyromagnetic ratio, moments of inertia of odd nuclei, moments of inertia in excited states, comparison with the experiment. Besides, the transition to an odd proton reduces the moment of inertia much less than the transition to an odd neutron. The author thanks L. D. Landau, D. F. Zaretskiy and A. I. Larkin for interesting discussions as well as I. M. Pavlichenkov and M. G. Urin for the compilation of the tables. There are 1 figure, 3 tables, and 7 references, 3 of which are Soviet.

SUBMITTED:

February 13, 1959

Card 2/2

sov/56-37-1-38/64 21 (1). Migdal, A. B. Superfluidity and Moments of Inertia of Nuclei (Sverkhtekuchest' AUTHOR: i momenty inertsii yader) TITLE: Zhurnal eksperimental noy i teoreticheskoy fiziki, 1959, Vol 37, PERIODICAL: Nr 1(7), pp 249 - 263 (USSR) The author discusses a method of investigating the superfluidity in systems of finite dimensions. The moments of inertia are ABSTRACT: calculated by this method in quasi-classical approximation which leads to a satisfactory agreement with the observed values of the moments of inertia. Also the calculated increase of the moment of inertia in the transition from an even-even to an even-odd nucleus, and also the gyromagnetic ratio for rotating nuclei, agree with the experiment. Thus, these results confirm the hypothesis of superfluidity of nuclear matter. Besides, the superfluidity of nuclear matter may lead to interesting macroscopic phenomena if there are stars with a core of neutrons. Such a star would be in the state of superfluidity, and the transition temperature would correspond to the value of 1 Mev. The paper is divided as follows: Introduction, a method of investigating a system of finite dimensions, Card 1/2

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The Interaction Between Electrons and Lattice Vibrations SOV/56-34-6-10/51 in a Normal Metal

> quasiparticle lies sufficiently near to the Fermi surface thedamping results from the interaction between the electrons. This interaction is caused by the exchange of phonons. The interaction between the electrons implies a damping that is proportional to the square of the "nearness" to the Fermi surface. The energy spectrum of the electrons is determined by the poles of the Green (Grin) function. There are 2 figures and 3 references, 2 of which are Soviet.

ASSOCIATION: Moskovskiy inzhenerno-fizicheskiy institut (Moscow Engineering-

Physics Institute)

SUBMITTED:

July 12, 1957(initially) and March 20, 1958 (after revision)

Card 3/3

The Interaction Between Electrons and Lattice Vibrations SOV/56-34-6-10/51

 $\lambda_{\mathbf{a}}$  is small. M denotes the mass of the ion and  $\lambda_{\mathbf{a}}$  a dimensionless parameter which does not contain the masses

of the ions. The author first introduces the propagation functions G and D of the electron and phonon :

G =  $i\langle \text{TW}(;) \text{W}^{4}(2) \rangle$ , D =  $i\langle \text{T}\phi(1) | \phi(2) \rangle$ . The average is determined with respect to the ground state of the system. The Dyson (Dayson) equations connect D and G with the vertex part P. Then a formula for the interaction energy is given. The next part of this paper deals with the vertex

part. It is different from  $\Gamma_p = 1$  by a quantity with the order of magnitude  $M^{-1/2}$ . The author discusses the first order correction of  $\Gamma$  according to the disturbation theory. The calculations are discussed step by step. Finally one obtains an expression of the type  $\Gamma=1+O(M^{-1/2})$ . This result is not changed by taking into account the graphs of the higher orders. The following parts of this paper deal with the Green (Grin) functions of the phonon and of the electron. The damping of the electron excitations results from the emission of phonons. If the energy of the

Card 2/3

in a Normal Metal

AUTHOR:

Migdal, A. B.

sov/56-34-6-10/51

TITLE:

The Interaction Between Electrons and Lattice Vibrations in a Normal Metal (Vzaimodeystviya elektronov s kolebaniyami

reshetki v normal'nom metalle)

PERIODICAL:

Zhurnal eksperimental noy i teoreticheskoy fiziki, 1958,

Vol 34, Nr 6, pp 1438 - 1446 (USSR)

ABSTRACT:

With the method developed in this paper one may - in the case of a normal metal - take into account the interaction of the electrons with the lattice vibrations. In this method it is not necessary to assume that the interaction is weak. In the case of superconduction the necessary condition for the applicability of the usual methods of field theory is not valid. Therefore the extension of the method (developed in this paper for normal metals) as regards the case of superconductivity requests a separate investigation. The energy spectrum of the Hamiltonian given by H. Fröhlich (Ref 2) cannot be obtained with the methods of disturbation theory. With the methods of field theory the energy spectrum can be found as a series with respect to the powers of M-1/2 without assuming that

Card 1/3

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An Application of Quantum Field Theory Methods to the Many- 56-1-22/55 -Body Problem.

arbitrary weak electromagnetic field are investigated. There are 3 figures and 8 references, 5 of which are Slavic.

ASSOCIATION: Moscow Engineering and Physical Institute (Moskovskiy institut).

SUBMITTED: July 12, 1957 (initially) and October 24, 1957 (after revision).

AVAILABLE: Library of Congress

Card 3/3

APPROVED FOR RELEASE: 06/23/11: CIA-RDP86-00513R001033800004-6

56-1-22/55

The kernel of the phonon determines the energy and the damping of the excitations of the lattice. At first the kernel G(pE) is written down for one particle, and then the author passes over to a Fourier representation. Subsequently, the properties of the kernel in the complex plane are investigated, and the interrelation of the kernel of one particle with the spectrum of the excitations is determined. The behaviour of the kernel at great positive times is also studied. The energy and the damping of the excitations are determined in the lower half plane by means of the real and imaginary part of the poles of the analytical propagation of G(pg). The kernel for one particle also permits the determination of other charachteristics of the system, e.g. the distribution of the particles on the different momenta. For the purpose of studying the energy spectrum and the behaviour of the system in weak external fields, it is necessary to investigate the kernel for two particles. This kernel for two particles is written down here explicitly, it is suited, for example, for studying the excited states of a system of N particles containing one particle and one hole. The case of forces of short range and the behaviour of a system in an

Card 2/3

MIGPAL A B

ALTHORS:

Galitskiy, V. M., Migdal, A. B.

56-1-22/56

TITLE:

An Application of Quantum Field Theory Methods to the Many-Body Problem (Primeneniye metodov kvantovoy teorii polya k zadache mnogikh tel).

PERIODICAL:

Zhurnal Eksperimental'noy i Teoreticheskoy Fiziki, 1958, Vol. 34, Nr 1, pp. 139-150 (USSR).

ABSTRACT :

In the present paper shows, that the energy and the damping of the quasi-particles depends on the poles of the dissipation function of a particle. The author here investigates a homogeneous unbounded system, wherein the momentum operator commutes with the Hamiltonian. In all Fermi-systems there obviously exist excitations analoguous to those in an ideal Fermi gas. It is convenient to study the properties of the excitations by means of the methods of the quantum theory of fields, by introducing the kernels of the system into the investigations. Apart from the kernels of the particles it is also possible to introduce the functions of the dissipation of the interaction between the particles, e.g. the kernel of the phonon represents this dissipation function in the problem of electrons in a metal being in interaction with the lattice.

Card 1/ 3

Plasma Physics and the Problem (Cont.) SOV/1242

Belyayev, S.T., and G.I. Budker. Relativistic Plasma in Variable Fields 283

Budker, G.I. and S.T. Belyayev. Kinetic Equation for an Electronic Gas for Rare Collisions 330

AVAILABLE: Library of Congress (QC794.A38)

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## SOV/1242 Plasma Physics and the Problem (Cont.) TABLE OF CONTENTS: Braginskiy, S.I., and V.D. Shafranov. Plasma Filament With 3 Heat Loss to the Electrodes Braginskiy, S.I., and A.B. Migdal. Processes in a Plasma Column During Rapid Increase of a Current 20 Braginskiy, S.I., and V.D. Shafranov. Plasma Filament in 26 the Presence of a Longitudinal Magnetic Field Artsimovich, L.A. Passage of Large Currents Through Plasma in 81 the Presence of a Longitudinal Magnetic Field 87 Arsimovich, L.A. Magnetic Flux in a Compressing Cylinder Artsimovich, L.A. Analysis of the Equation of Compression of a Filament in the Presence of an External Magnetic Field 101 Card 2/5

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MIGDAL, AB.

21(7)

PHASE I BOOK EXPLOITATION SOV/1842

Akademiya nauk SSSR. Institut atomnoy energii

Fizika plazmy i problema upravlyayemkh termoyadernykh reaktsiy, t. II. (Plasma Physics and the Problem of Controlled Thermonuclear Reactions, t. 2) [Moscow] Izd-vo AN SSSR, 1958. 355 p. 3,000 copies printed.

Resp. Ed.: Leontovich, M.A., Academician.

PURPOSE: This collection contains previously unpublished work of members of the Institut atomnoy energii (Institute of Atomic Energy) of the Academy of Sciences of the USSR. It is intended for scientists interested in this field.

COVERAGE: This book is the second of four volumes of previously unpublished work of members of the Institute of Atomic Energy during the period 1951-58. The exploitation cards on the other volumes in this series have been released under the numbers 1241,1243, and 1244.

Card 1/5

APPROVED FOR RELEASE; 06/23/11: CIA-RDP86-00513R001033800004-6

KOGAN, V. I. and MIGDAL, A. B.

"The Electron Temperature Dependence of the Spectrum of the Bremsstrahlung of a Plasma." (Work - 1951); (and reworked in Preparation for publication); pp. 172-177.

"The Physics of Plasmas; Problems of Controlled Thermonuclear Reactions;" Vol. 1. 1958, published by Inst. Atomic Energy, Acad. Sci. USSR. resp. ed. M. A. Leontovich, editorial work V. I. Kogan.

Available in Library.

APPROVED FOR RELEASE: 06/23/11: CIA-RDP86-00513R001033800004-6

GALITSKIY, V. M. and MIGDAL, A. B.

"Dielectric Constant of a High Temperature Magnetized Plasma and the Evaluation of the Radiant Heat Conductivity." (Work - 1951); pp. 161-171.

"The Physics of Plasmas; Problems of Controlled Thermonuclear Reactions." Vol. 1. 1958, published by Inst. Atomic Energy, Acad. Sci. USSR. resp. ed. M. A. Leontovinh, editorial work V. I. Kogan.

Available in Library.

AUTHOR TITLE

The Bremsstrahlung and Forming of Pairs at High Energies in a Condemsed

56-4-1/52

Medium.

(Tormoznoye izlucheniye i obrazovaniye par v kondensirovannoy srede -

Russian)

PERIORICAL

Zhurnal Eksperim,i Teoret.Fiziki,1957,Vol 32,Nr 4, pp 633-646 (U.S.S.R.)

Received 7/1957

ABSTRACT

The present paper investigates the influence exercised by multiple scattering upon bremsstrahlung and forming of pairs. Formulae are given for the probability of bremsstrahlung and formation of pairs on the unit of the way in a condensed medium without restriction of the energies of electrons and quanta. For this purpose a connection between the transition probability and the matrix of density is set up. The probability of bremsstrahlung or formation of pairs must be averaged over all possible locations of the atoms of the scattering medium. Averaging of the number of transitions in the time unit is reduced to the search for the averaged matrix of density and to the determination of two formulae given here. Next, the author derives an equation for the averaged matrix of density, computation is followed step by step. Summation over the spin of the electron and polarization of the quantum in the formulae derived here can be carried out in the usual manner. In one of the formulae summation can be reduced to the determination of the trace of double-rowed matrizes. A formula is now given for the probability of bremsstrahlung, i.e. for the radiation of a quantum with an energy between k and k + dk per unit of way. An expression for the angular distribution of the quanta is also given. This angular distribution

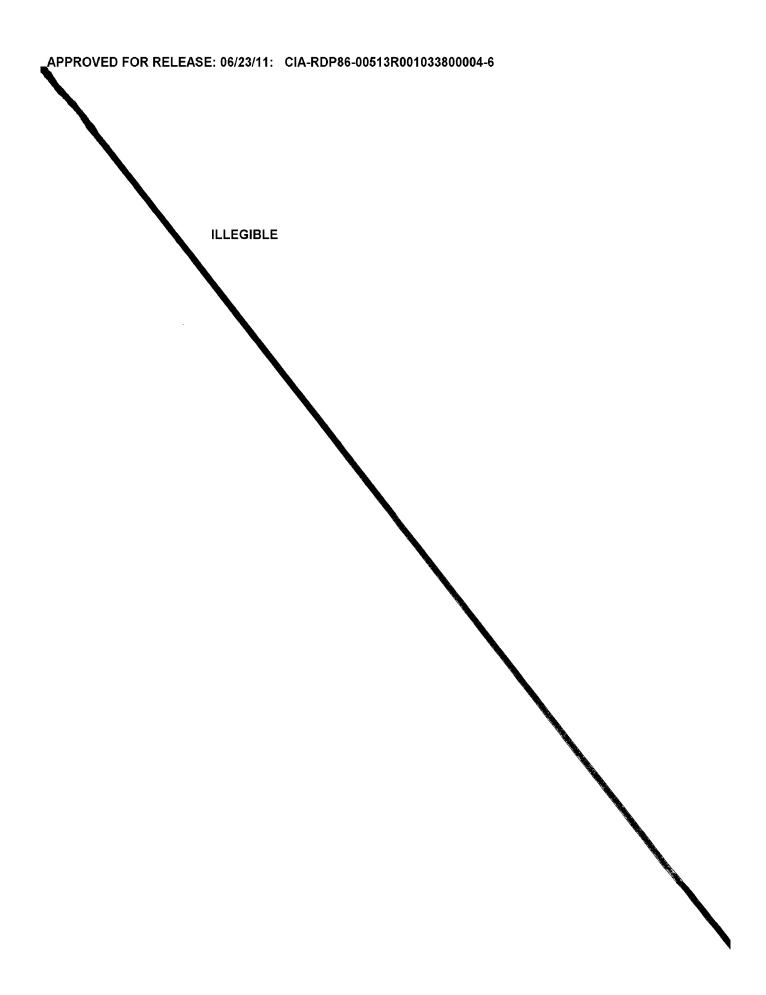
Card 1/2

PA - 2709 On the Distribution of FERMI Particles Which are in Interaction over the Momenta. se of GREEN'S function of a particle:

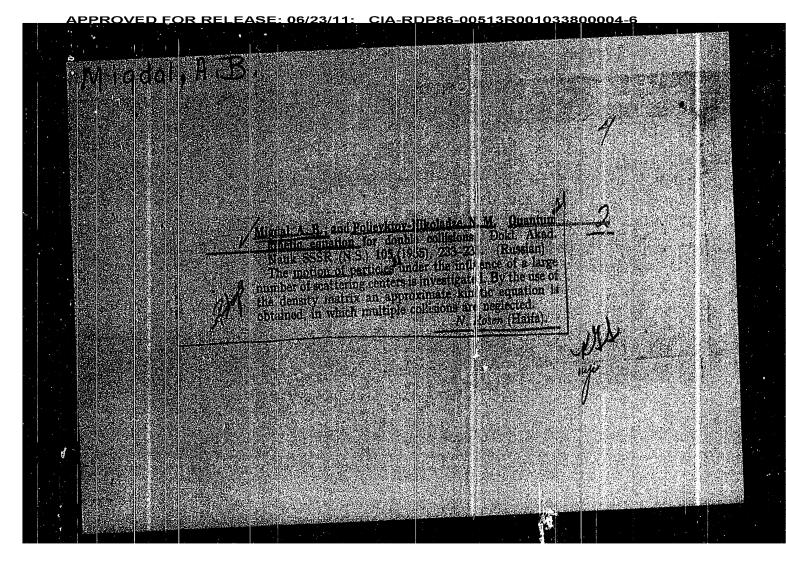
G (r,t,r,t)=i Teint \(r,t)=iH(t,-t,a)\(\psi + (r,a)=t,a)\(\psi + (r,a)=t,a)\(\psi + (r,a)=t,a)=t,a This GREEN'S function is then developed into a FOURIER series and brought into connection with the momentum distribution of the particles in the ground state. The following expression is obtained for this distribution:  $n(p) = i \int_{C} G(p, \xi) d\xi/2\pi$ . GREEN'S function must have poles which correspon to the quasiparticles. If p is near po, G(p, E) =Z/(  $E_p$ - E-i $\gamma$ (p))+ + f(p, E)applies. Here the function f(p, E) describes the damping of the quasiparticles and changes its sign at  $p = p_0$ . Z here denotes the renormalization constant of GREEN'S function. ASSOCIATION PRESENTED BY 22.11.1956 SUBMITTED Library of Congress AVAILABLE Card 2/2

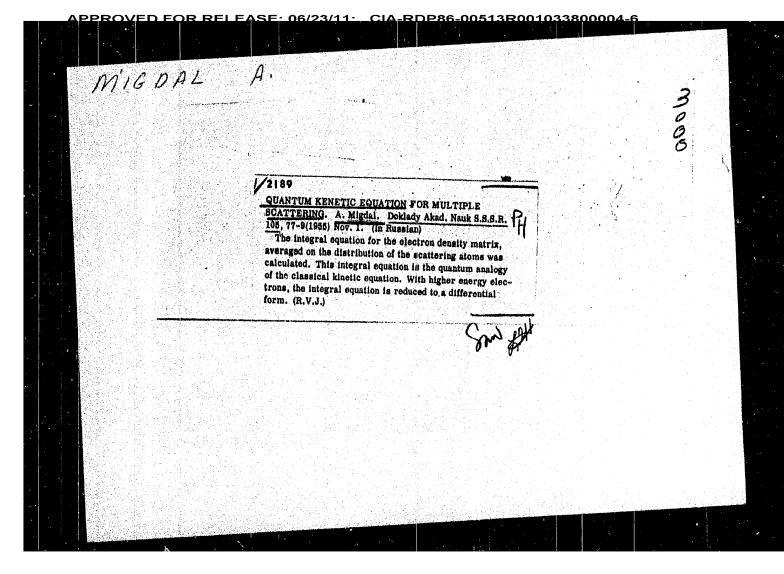
PA - 2709 On the Distribution of FERMI Particles Which Are in Interaction over AUTHOR (O raspredelenii vzaimodeystvuyushchikh Fermi-chastits po impulsam, TITLE Zhurnal Eksperim.i Teoret.Fiziki, 1957, Vol 32, Nr 2, pp 399-400 (U.S.S.R.) PERIODICAL The system investigated here consists of a great number of particles Received 5/1957 which are in interaction. The excited states of the system should contain states the energy of which may be represented as a sum of the energies ABSTRACT of the quasiparticles. A formula is given for the energy of the quasiparticles. The quasiparticles have a damping that is proportional to (p - po)2. (Here p denotes the momentum of the quasiparticle, po - the maximum momentum of the FERMI-filling for the quaiparticles, vo= v(po) the velocity of the quaiparticles on the FERMI surface. p>po corresponds to a quasiparticle and p po to a hole). In the case of sufficiently strong interaction, an excited state of the system for p which are near  $p_0$ , cannot therefore be described by means of quadparticles. At  $p \rightarrow p_0$ the quasiparticles are able to describe the state of the system also in the case of strong interaction. The author here shows that the distribution of the particles over the momenta at the basic state hasajumb at p = po, and this applies to any interaction. Here, it must be born in mind that the distribution of the particles (and not of the quasiparticles) over the momenta is concerned. The following equation applies in the ca-Card 1/2

CIA-RDP86-00513R001033800004



MIGDAL, A. B., GALITSKIY, V. M. and LANDAU, L. D. "The Disintegration of the Deuteron by the Coulomb Field of the Nucleus" a paper Presented at the International Conference on Nuclear Reactions, Amsterdam, 2-7 July 1956. D551274





"Theory of Scattering in a Guard-Glassical Ambreviration," Zg. 1997. 1877. 1878. 1989. 1989. 1989. MICHAL, A. B. The wave-function is represented as \$ = \$ \frac{1}{2} \text{ is the review is then \$4 \frac{1}{2} \text{ is then \$4 \frac{1}{2} \text{ is the review is the re

LEVE RELEASE: 06/23/11: CIA-RDP86-00513R001033800004-6

MIGDAL, A.B.
USSR/Nuclear Physics - Meson Formation

FD-1882

Card 1/1

Pub. 146-2/21

Author

: Migdal, A. B.

Title

: Meson formation at an energy close to the threshold

Periodical:

Zhur. eksp. i teor. fiz. 28, 10-12, 1955

Abstract

: The energy spectrum of motions originating in collisions of two nucleons is determined. The highest probability is attributed to formation of nucleons of weak mutual energy. Therefore the most probable meson energy should approach the maximum meson energy, in accordance with the conservation law. The correlation between the cross section of deuteron formation in reactions  $p + p = n + p + \pi^+$ ;  $n + p = p + n + \pi^0$ ;  $n + n = n + p + \pi^-$  and cross sections of formation of free neutrons and protons with parallel spins is obtained. Three USSR and two US references.

. Institution:

Submitted : February 23, 1954

PPROVED FOR RELEASE: 06/23/11: CIA-RDP86-00513R001033800004-6

MIGDAL, A. B. USSR/Nuclear Physics - Reactions

FD-1881

Card 1/1

Pub. 146-1/21

Author

: Migdal, A. B.

Title

Theory of nuclear reactions producing slow particles

Periodical: Zhur. eksp. i teor. fiz. 20, 3-9, 1955

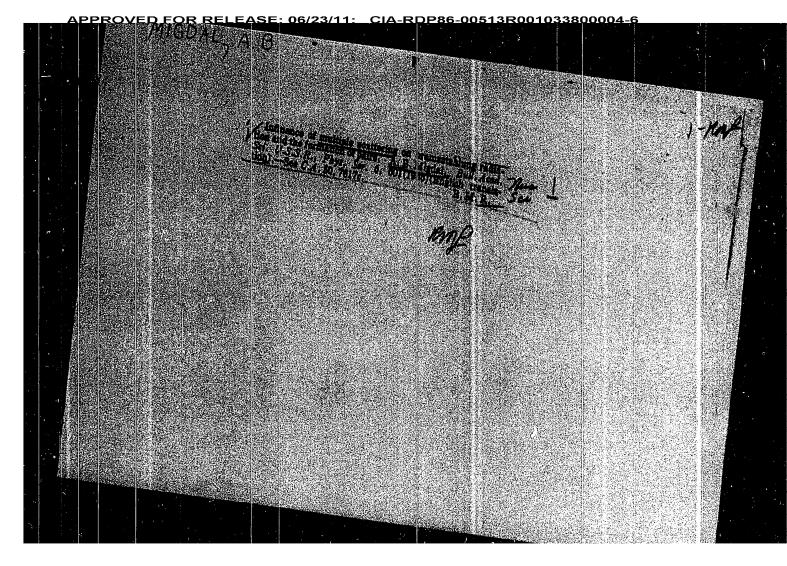
Abstract

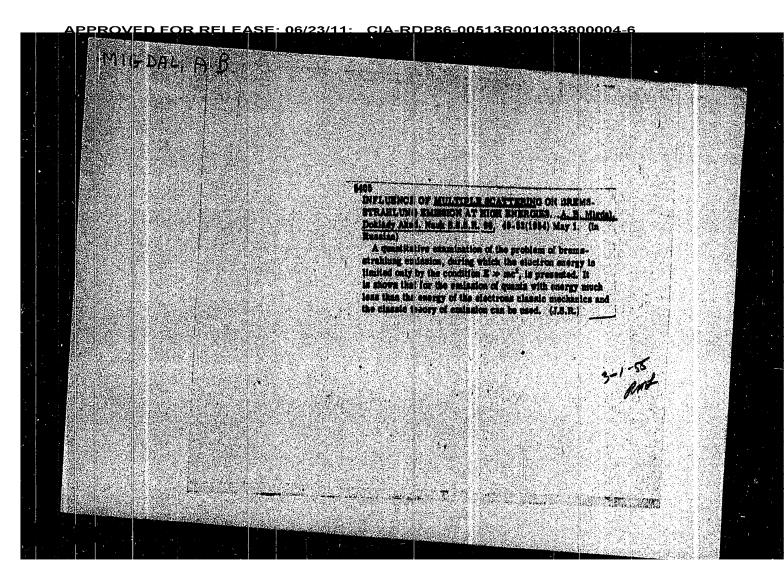
: The energy and angular distribution of slow nucleons produced as result of nuclear reactions are analyzed. The probability of nucleon formation in a bound state is estimated. Author thanks B. T. Geylikman, I. Ya. Pomeranchuk and Ya. A. Smorodinskiy for discussions. One USSR and two US references.

Institution:

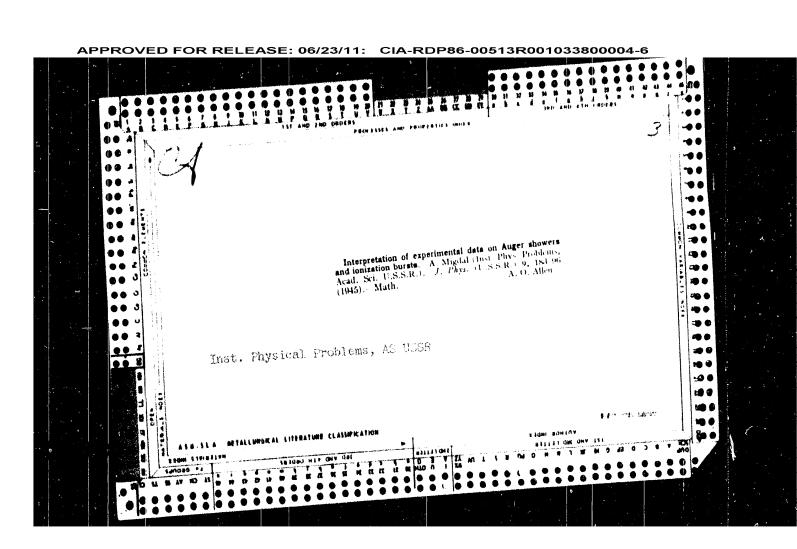
Submitted: February 23, 1953

MIGDAL, A.B. Effect of multiple scattering on bremsstrahlung and on pair formation. Izv.AN SSSR.Ser.fiz.19 no.6:665 N-D \*55.(MLRA 9:4) 1. Akademiya nauk SSSR. (Cosmic rays) (Nuclear physics)





168761 MIGDAL, A. B. Jun 50 USSR/Nuclear Physics - Mesons "Artificial 17-Mesons," A. B. Migdal, Ya. A. Smorodinskiy "Uspekh Fiz Nauk" Vol XLI, No 2, pp 133-153 Discusses artificial production of  $\pi$ -mesons using Berkely cyclotron in 1947. Lists only four Soviet sources, all for the introduction on varitrons. 168761

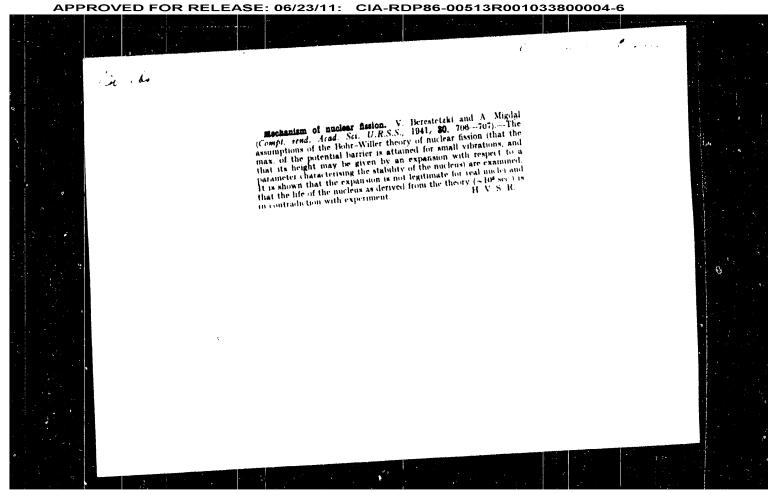


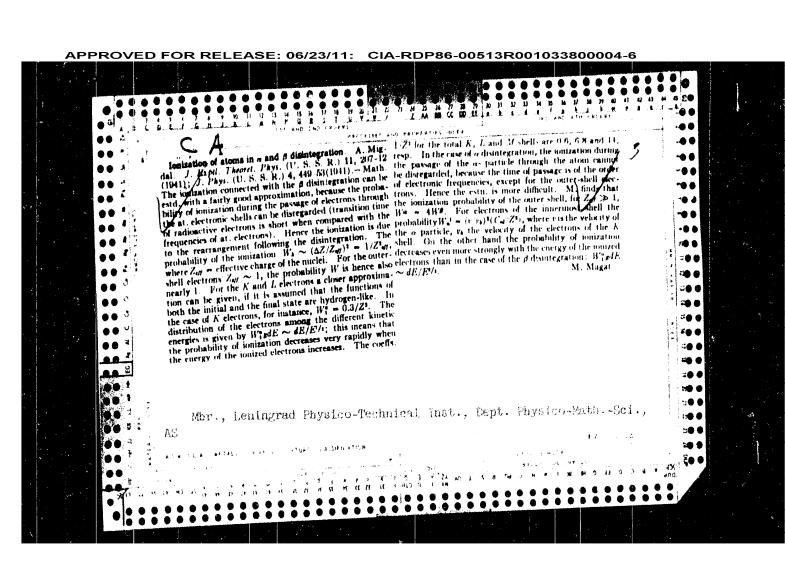
and Hoffman's Impacts," MIGDAL, A. 3. "Analysis of Some Experimental Late on Appur Chowers," sime ortion contests of consists of 1974 in 1984. IAM-Ser Fiz, Vol 2, No 3, 1945

APPROVED FOR RELEASE: 06/23/11: CIA-RDP86-00513R001033800004-6 • • • • Equilibrium spectrum of the soft component of high-energy cosmic rays. A. Migdal (Inst. Phys. Problems, Acad. Sci. U.S.S.R.). J. Phys. (U.S.S.R.) 9, 87-92 (1945); cf. C.A. 39, 32014. Theoretical. The equi-spectrum of the soft component in air is calcolar numerically for energies several times greater than the crit. energy. E. J. Rosenbaum METALLUFGICAL LIFEHATURE (LASSIFICATION # 4 # j ....

## APPROVED FOR RELEASE: 06/23/11: CIA-RDP86-00513R001033800004-6 Pair creation in nuclear collisions. A. Migdal (Inst. Phys. Problems. Acad. Sci. U.S.S.R.). J. Phys. (U.S.S.R.) 9, No. 1, 45-51 (1945). The probability of pair screation is calcid, where one of the colliding particles is charged. The probability in a dreay and fission of U.S. calcid, by expanding the field in pendergriphs and point of the probability in a dreay and fission of U.S. calcid, by expanding the field in pendergriphs and pendergriphs. Inst. for Phys. Problems, AS USSR ASH-SLA METALLURGHCAL LITERATURE CLASSIFICATION

## APPROVED FOR RELEASE: 06/23/11: CIA-RDP86-00513R001033800004-6 PROCESSES AND PROPERTIES INDEX Quadrupole and dipole $\gamma$ -radiation of nuclei. A. Migdal, J. Phys. (U. S. S. R.) 8, 331–6(1944). Calcus, are made of the magnitude of matrix elements for dipole and quadrupole transitions. These calcus, are used to explain the fact that the intensities of corresponding $\gamma$ -rays from the two transitions with energies about 1 in e.v. are of the same order of magnitude. For energies much greater than 1 in e.v. the intensities of dipole transition must be much greater than the intensities of quadrupole transitions. Institute of Physical Problems. gig mit my pagetie.





MIGDAL, A. B.]

"The Mechanism of the Fission of Heavy Nuclei," J. Phys (USSR), Vol h, p.283, 19h1

W-38h, 16 Apr h8

Also: Dok. AN 30, No. 8, 19h1.

## APPROVED FOR RELEASE: 06/23/11: CIA-RDP86-00513R001033800004-6 BC Pads of the mesotron tracks observed in an expansion chamber. A. Migdal and J. Pomerantschula (Compl. rend. Acad. Sci. If R.S.S., 1940, 37, 652–653)....The absence of electron tracks at the ends of mesotron tracks in cloud-chamber photographs may be explained by diffusion of the mesotron over a considerable distance after its energy has decreased too far for it to produce appreciable ionisation, and before disintegration occurs. The distance covered in this diffusion region is calc. as ~1.0 cm. for E = 10 ky, in air

### APPROVED FOR RELEASE: 06/23/11: CIA-RDP86-00513R001033800004-6 Scattering of neutrons in paramagnetic substances. A. B. Migdal -J Expli Theoret, Phys. (U. S. S. R.) 10, 50-44(1010), cf. ( i. 13, 1688). The scattering cross section of thermal our consumparamagnetic substances was called on the basis of wave mechanics with special emphasis on this case of accounter. In this case the cross section is given by the expression $\sigma = 0.05$ ( $\sigma = 0.05$ ) of the content, where C is about 1 and B is the magneton notice T (T T T) of the element in question. For Ce. Sin. cm., where C is about a and B is the magnetian $(g \times T(1+T))$ of the element in question. For Cc, Sin, Dy and Fr, B is (2.51, 0.84, 10.6) and 9.60, resp. In the case of Fu $B \neq 0$ and the scattering must be about in spite of the presence of paramagneticm. Rok select G most Leningred Physico-Tech. Inst., Leningred

#### APPROVED FOR RELEASE: 06/23/11: CIA-RDP86-00513R001033800004-6 (0) lonization of atoms and pair creation in the course of nuclear reactions. A. B. Migdal. Bull. acad. w. U. R. S. S., Ser. phys. 4, 287-80m English, 2809 (1940); cf. C. J. 35, 25%. The transition probabilities for the unterationic processes in which the disturbance period is short as compared with proper period of electrons can be calcillated expanding the original wave function of the system, over the wave functions corresponding to the new state. By this method the formulas have been obtained for the probability of pair formation in a proton-healtron collision (for relativistic velocities of colliding particles), and also for the probability of pair formation accompanying the fission of the atomic nucleus. Roksalana Gamow Physico-Technical Institute of the AS MaSR, Lamingrad. BETALLUNGKIAL SITERATURE CLASSIFICATION

# Scattering of neutrons in terromagnetic A Minds (completed and the control and

L 45093-66 EWT(m)/T

ACC NR: AP6024873 SOURCE CODE: UR/0056/66/051/001/0135/0146

AUTHOR: Migdal, A. A.; Polyakov, A. M.

ORG: none

TITLE: Spontaneous violation of strong interaction symmetry and the absence of zero-mass particles

SOURCE: Zhurnal eksperimental'noy i teoreticheskoy fiziki, v. 51, no. 1, 1966, 135-146

TOPIC TAGS: vector meson, Feynman diagram, quantum electrodynamics, C invariance, PARTICLE, INTERACTION

ABSTRACT: The existence of zero mass particles in the presence of spontaneous violation of symmetry is considered. Summation of all Feynman diagrams yields an equation which is identical to the Bethe—Saltpeter equation for the wave function of a zero mass scalar bound state (zeron) in the annihilation channel ab of the particles for the difference between the mass of erators  $M_{a}(p)$ —M (p) of particles a snd b in a supermultiplet. It is shown that in spontaneous violation of symmetry in a Yang-Mills type theory with vector mesons, the zerons interact only with virtual particles and hence are unobservable. On the other hand, vector mesons acquire a mass despite the generalized guage invariance. It is proved that an asymmetric solution corresponds to a minimal vacuum energy and that a consequence of C-invariance of the solution is the conservation of strangeness. Orig. art. has: 26 formulas.

SUB CODE: 20/ SUBM DATE: 30Nov65/ ORIG REF: 003/ OTH REF: 010

APPROVED FOR RELEASE: 06/23/11: CIA-RDP86-00513R001033800004-6

ACCESSION NR: AP4012547

In such a potential, it is possible to apply the results obtained to a calculation of the boundary energy of Fermi nuclei. The final formula can be represented in the form

 $e_0/e_\infty = 1 + A^{-1/1}/(N/A)$ ,

where f(x) is given by

x = 0.50 0.52 0.54 0.58 0.58 0.00 0.62 0.64 0.66 0.68 /(x) = 1.74 1.71 1.51 1.50 1.40 1.34 1.20 1.13 1.00 0.90

and which is accurate to within 1--2 MeV. Orig. art. has: 26 formulas.

ASSOCIATION: Moskovskiy fiziki-tekhnicheskiy institut (Moscow Physicotechnical Institute)

Card 2/32

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ACCESSION NR: AP4012547

S/0056/64/046/001/0213/0217

AUTHORS: Gurvits, S. A.; Migdal, A. A.; Polyakov, A. M.

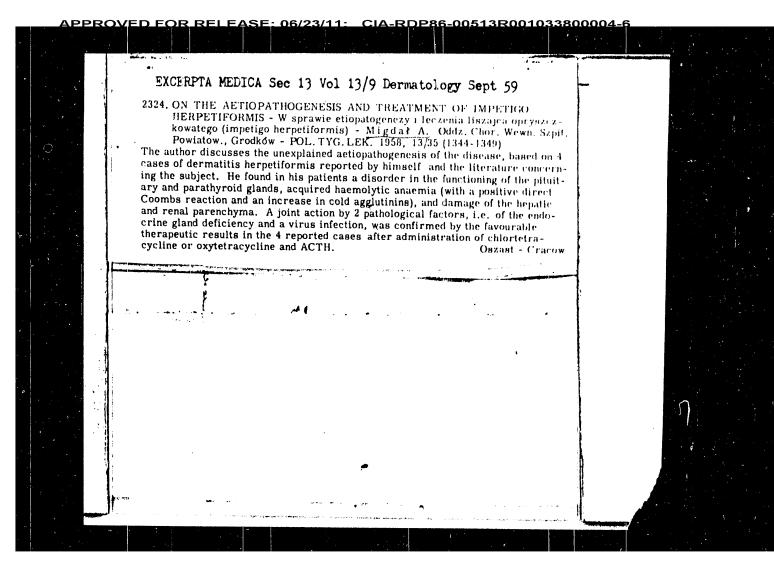
TITLE: Boundary energy of a Fermi gas in a potential well

SQURCE: Zhurnal eksper. i teoret. fiz., v. 46, no. 1, 1964, 213-217

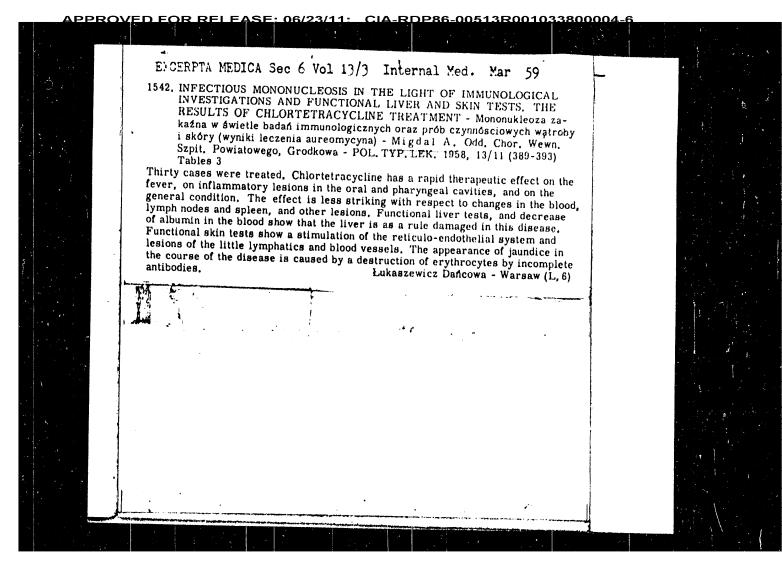
TOPIC TAGS: Fermi gas, quantum mechanics, potential well, quantization, Fermi energy, boundary Fermi energy, neutron Fermi energy, Beavy nucleus Fermi energy, Fermi nucleus

ABSTRACT: A quasi-classical quantization condition is obtained for a spherically symmetrical potential and is used to obtain the first two terms of an expansion of the number of particles, expressed in the form of a function of the Fermi-gas boundary energy, in powers of the dimensions of the system for this potential. The method given makes it possible to make similar calculations for any potential well with a diffused edge. By regarding nucleons as a Fermi gas

Card 1/37



-RDP86-00513R001033800004 MIGDAL, Adam (Grodkow, Powiatowa Przychodnia Skorno-Weneryczna) Paravaccinia & vaccination vaccinia. Polski tygod. lek. 13 no.31: 1199-1203 4 Aug 58. 1. Z Pow. Przychodni Skorno-Wenerycznej w Grodkowie; kierownik: dr A. Migdal; z Pow. Inboratorium Analit. Bakteriol. w Grodkowie; kierownik: dr A. Migdal. (VACCINIA paravaccinia & vaccinia (Pol))



APPROVED FOR RELEASE: 06/23/11: CIA-RDP86-00513R001033800004-6 MIGDAL, Adam, (Grodkow, pow. Opole, Szpital Powiatowy.) March Control of the Control Biochemical disorders of the organism during pellagroid. Przegl. derm., Warsz. 7 no.1:49-58 Jan-Feb 57. 1. Z Oddzialu Chorob Wewnetrznych Szpitala Powiatowego w Grodkowie Ordynator: dr A. Migdal. Z Powiatowego Iaboratorium Analitycznego w Grodkowie. Kierownik: dr A. Migdal. (PELLAGRA biochem. disord. in pellagroid (Pol))

CIA-RDP86-00513R001033800004-6 APPROVED FOR RELEASE: 06/23/11: MIGDAL, Adam, (Grodkow) Azotniak and its relation to health of workers during spreading. Polski tygod. lek. 11 no.11:489-494 12 Mar 56. (FERTILIZERS, injurious effects, (Pol))

MIGDAL, Adam, Grodkow, ul. Sienkiewicza 56

Behavior of eosinophilic leukocytes and of skin and liver function tests in allergic skin diseases treated with nitrogen mustard in small doses. Polski tygod. lek. 10 no.1:7-13 3 Jan 55

Z Pow. poradni skorno-wener. w Grodkowie; kier. dr. med. A.Migdal (SKIN, diseases allergic, eff. of nitrogen mustard ther. on leukocytes behavior & skin & liver funct. test)
 (NITROGHN MUSTARDS, effects on leukocytes behavior, skin & liver funct. test in ther. of allergic skin dis.)

(SKIN, physiology
eff. of nitrogen mustards in ther. of allergic skin dis.)
(LIVER FUNCTION TEST, in various diseases
skin dis., allergic, eff. of nitrogen mustards ther.)

#### MIGDAL, Adam

Nitrogranulogen therapy of allergic diseases of the gastrointestinal system. Polski tygoi. lek. 9 no.12:362-366 22 Mar 54.

1. Z Odds. Chorob Wawnetrznych Szpitala Pow w Grodkowie, zast, dyrektora: dr. A.Migdal.

(ALLERGY, manifestations,

gastrointestinal, there, nitrogen mustards)

(GASTROINTESTINAL DISEASES,

allergy, ther., nitrogen mustards)

(NITROGEN MUSTARDS, therapeutic use,

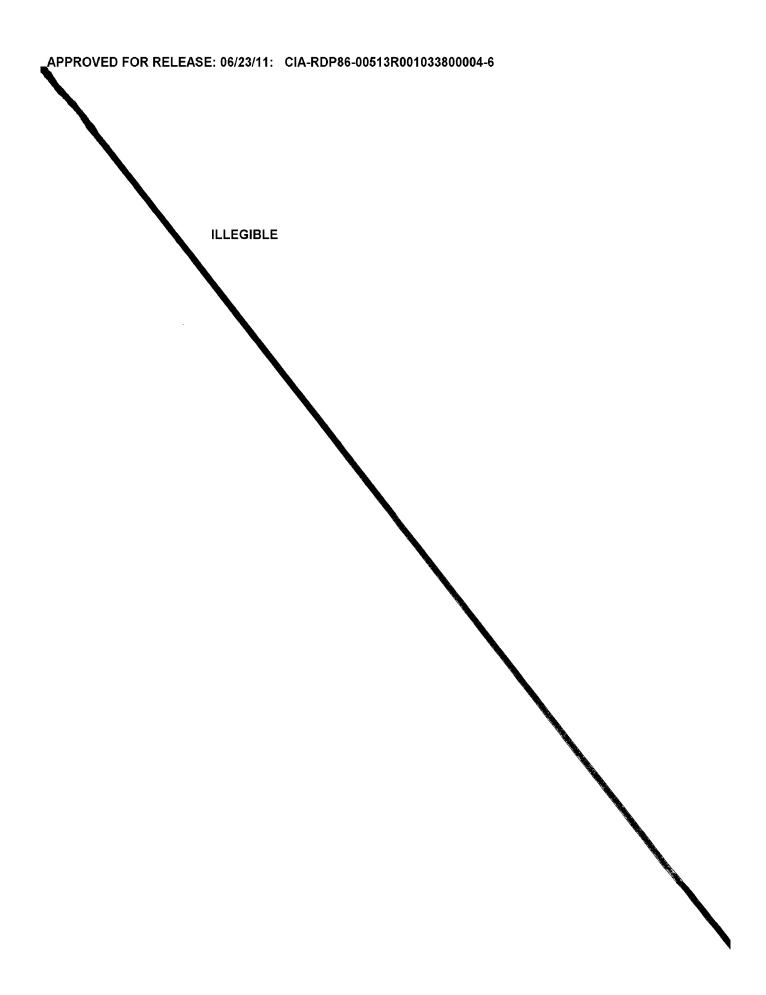
allergy of gastrointestinal system)

MIGDAL, A.; KALUZYNSKI, H.

Attempts to treat allergic skin diseases by nitrogranulogen. Polski tygod. lek 7 no. 41:1277-1281 13 Oct 1952. (CLML 24:1)

1. Of Grodkow District Dermato-Venereological Consultation Center (Head--Adam Migdal, M.D.) and of Niemodlin District Hospital (Director--Henryk Kaluzynski, M.D.).

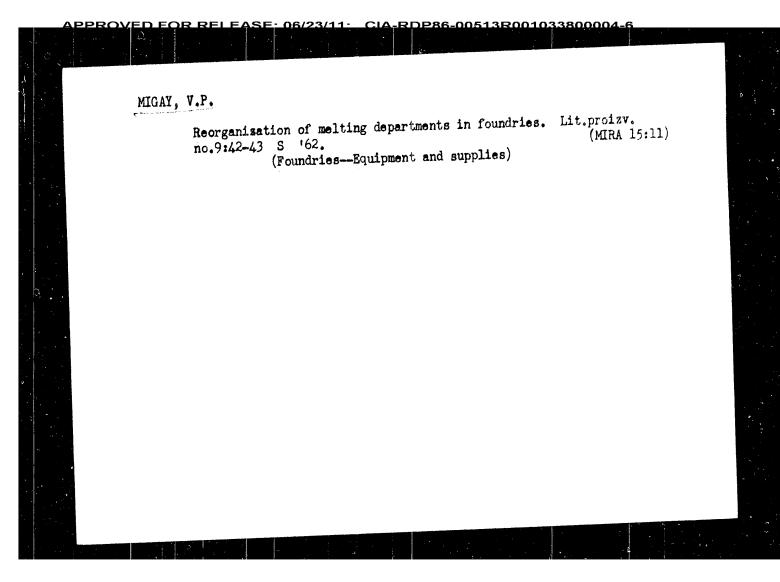
<u> APPROVED FOR RELEASE: 06/23/11: CIA-RDP86-00513R001033800004-6</u> MIGDA, Tadeusz, mgr., inz. Pipe rolling mills. Przegl techn 81 no.18:17-19 '60.



MIGAY, V.P. Effect of gases on the structure and the graphitization of cast iron. Metalloved. i term. obr. met. no.11:35-36 N 163. (MIRA 16:11) 1. Vsesoyuznyy zaochnyy mashinostroitel'nyy institut.

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CIA-RDP86-00513R001033800004-6 MIGAY, V.P. Effect of vacuuming on the structure and graphitization of cast iron. Lit. proizv. no.2:20-21 F '63. (MIRA (Cast iron--Metallography) (Vacuum metallurgy) (MIRA 16:3) MIGAY, V.P. Vacuum treatment of gray cast iron. Lit. proizv. no.1:25-26
Ja 163. (MIRA 16:3) (Cast iron)
(Vacuum metallurgy)



MIGAY, V.P. Reducing the gas content of a foundry atmosphere during electric arc furnace operations. Lit.proizv. no.7:42 Jl '62. (MIRA 16:2) (Air--Purification)

<u> APPROVED FOR RELEASE: 06/23/11: CIA-RDP86-00513R001033800004-6</u> WINE, V.P. Mechanism for the operation of cupoda bottoms. Lit. proizv. no.12:31 D 161. (MIRA 14:12) (Cupola furnaces)

On the utilization of centrifuges in...

S/128/60/000/010/013/016/XX A033/A133

height of rotor rim - 120 mm; full rotor holding capacity - 240 liters; maximum rotor speed - 430 rpm; drive motor power - 40 kw; drive motor speed - 980 rpm; output (molding sand containing less than 2 - 3% clay) - 5,000 kg/hour; overall diemsnions: length (with motor) - 3,490 mm; height - 2,380 mm; width - 2,240 mm; weight (with motor) - 7,670 kg. There is 1 figure.

Card 3/3